

# SELECTED PROBLEMS IN PROBABILITY

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*Dedicated in friendship to Catriona Byrne*

ABSTRACT. This celebratory article contains a personal and idiosyncratic selection of a few open problems in discrete probability theory. These include well known questions concerning the Lorentz model and self-avoiding walks, and also some problems of percolation-type. The author hopes the reader will find something to leaven winter evenings, and perhaps even a project for the longer term.

## CONTENTS

Personal remarks	2
1. Introduction	2
2. Self-avoiding walks	2
2.1. Origins	2
2.2. Asymptotics	3
2.3. Self-avoiding walks in a random environment	4
3. Lorentz scatterers	4
3.1. Background	4
3.2. Ehrenfest wind/tree model	4
3.3. Poisson mirrors	5
3.4. Manhattan pinball	7
4. Two stochastic inequalities	7
4.1. Bunkbed inequality	7
4.2. Negative correlation	8
5. Randomly oriented square lattice	9
6. Enhancements	9
7. Two stochastic epidemics	11
8. Acknowledgements	12
References	12

## PERSONAL REMARKS

The editorial team of Springer Mathematics has become almost family for many of us worldwide, with Catriona Byrne at its heart. She has come to know us better than we know ourselves, always with sympathy, and with an honest and constructive approach to occasionally challenging areas of professional debate. Through our numerous collaborations, she and I have kindled a warm friendship that will persist into the next phase of our adventures. We wish her many happy years free from the woes of authors, editors, readers, and publishers.

## 1. INTRODUCTION

Probability has been a source of many tantalising problems over the centuries, of which the St Petersburg paradox, Fermat's problem of the points, and Bertrand's random triangles feature still in introductory courses. Whereas the conceptual problems of the past are now largely resolved, contemporary questions arise frequently where the intuitive apparatus of sub-fields collide. Many prominent problems are to be found at the conjunction of probability and discrete geometry. This short and idiosyncratic article summarises some of these. This account is personal and incomplete, and is to be viewed as a complete review of nothing. The bibliography is not intended to be complete, and apologies are extended to those whose work has been omitted.

The questions highlighted here vary from the intriguing to the profound. Whereas some may seem like puzzles with limited consequence, others will require new machinery and may have far-reaching implications.

## 2. SELF-AVOIDING WALKS

**2.1. Origins.** Self-avoiding walks were first introduced in the chemical theory of polymerisation (see [17, 46]), and their properties have received much attention since from mathematicians and physicists (see, for example, [7, 23, 42]).

A path in an infinite graph  $G = (V, E)$  is called *self-avoiding* if no vertex is visited more than once. Fix a vertex  $v \in V$ , and let  $\Sigma_n(v)$  be the set of  $n$ -step self-avoiding walks (SAWs) starting at  $v$ . The principal combinatorial problem is to determine how the cardinality  $\sigma_n(v) := |\Sigma_n(v)|$  grows as  $n \rightarrow \infty$ , and the complementary probability problem is to establish properties of the shape of a randomly selected member of  $\Sigma_n(v)$ . Progress has been striking but limited.

**2.2. Asymptotics.** It is now regarded as elementary that the so-called *connective constant*  $\kappa = \kappa(G)$ , given by

$$\log \kappa = \lim_{n \rightarrow \infty} \frac{1}{n} \log \sigma_n(v),$$

exists when  $G$  is quasi-transitive, and is independent of the choice of  $v$ . Thus, in this case

$$\sigma_n(v) = \kappa^{(1+o(1))n}.$$

The correction term is much harder to understand. We shall not make precise the concept of a  $d$ -dimensional lattice, but for definiteness the reader may concentrate on the hypercubic lattice  $\mathbb{Z}^d$ .

**Conjecture 2.1.** *For  $d \geq 2$  there exists a critical exponent  $\gamma = \gamma_d$  such that the following holds. Let  $G$  be a  $d$ -dimensional lattice. There exists a constant  $A > 0$  such that*

$$(2.1) \quad \sigma_n \sim An^{\gamma-1}\kappa(G)^n \quad \text{as } n \rightarrow \infty.$$

Furthermore,

$$\gamma = \begin{cases} \frac{11}{32} & \text{when } d = 2, \\ 1 & \text{when } d \geq 4. \end{cases}$$

A logarithmic correction is in fact expected in (2.1) when  $d = 4$ . See [7, 42] for further discussion and results so far, and the papers [29, 30] of Hara and Slade when  $G = \mathbb{Z}^d$  with  $d \geq 5$ , for which case they prove that  $\gamma = 1$ . Of particular interest is the case when  $G = \mathbb{H}$ , the hexagonal lattice. By a beautiful exact calculation that verifies an earlier conjecture of Nienhuis [45] based in conformal field theory, Duminil-Copin and Smirnov [15] proved that

$$\kappa(\mathbb{H}) = \sqrt{1 + \sqrt{2}}.$$

The proof reveals a discrete holomorphic function that is highly suggestive of a connection to a Schramm–Loewner evolution (see [34]), namely the following.

**Question 2.2.** *Does a uniformly distributed  $n$ -step SAW from the origin of  $\mathbb{H}$  converge weakly, when suitably rescaled, to the Schramm–Loewner random curve  $\text{SLE}_{8/3}$ ?*

Progress on this question should come hand-in-hand with a calculation of the associated critical exponent  $\gamma = \frac{11}{32}$ .

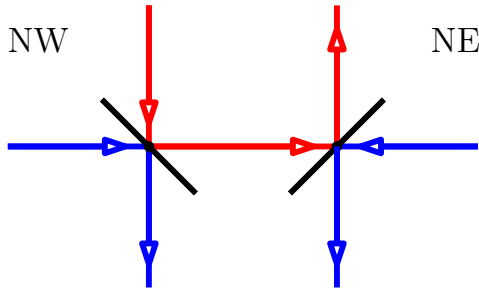


FIGURE 3.1. A NW and a NE mirror. Each is reflective on either side.

**2.3. Self-avoiding walks in a random environment.** How does the sequence  $(\sigma_n)$  behave when the underlying graph  $G$  is random? For concreteness, we consider here the infinite cluster  $I$  of bond percolation on  $\mathbb{Z}^2$  with edge-density  $p > \frac{1}{2}$  (see [20]).

**Question 2.3.** *Does the limit  $\mu(v) := \lim_{n \rightarrow \infty} \sigma_n(v)^{1/n}$  exist a.s., and satisfy  $\mu(v) = \mu(w)$  a.s. on the event  $\{v, w \in I\}$ ?*

Related discussion, including of the issue of deciding when  $\mu(v) = p\mu(\mathbb{Z}^2)$  a.s. on the event  $\{v \in I\}$ , may be found in papers of Lacoïn [35, 36]. The easier SAW problem on (deterministic) weighted graphs is considered in [25].

### 3. LORENTZ SCATTERERS

**3.1. Background.** The scattering problem of Lorentz [38] gives rise to the following general question. Scatterers are distributed randomly about  $\mathbb{R}^d$ . Light is shone from the origin in a given direction. Under what circumstances is the light ray: (i) bounded, (ii) unbounded, (iii) diffusive? While certain special cases are understood, the general question remains open. The problems mentioned here are concerned with *aperiodic* distributions of scatterers; the periodic case is rather easier.

**3.2. Ehrenfest wind/tree model.** The following notorious problem on the square lattice  $\mathbb{Z}^2$  has resisted solution for many years. Let  $p \in [0, 1]$ . At each vertex of  $\mathbb{Z}^2$  is placed a mirror with probability  $p$ , or alternatively nothing. Mirrors are plane and two-sided. Each mirror is designated a *north-east* (NE) mirror with probability  $\frac{1}{2}$ , or alternatively a *north-west* (NW) mirror. The states of different vertices are independent. The meanings of the mirrors are illustrated in Figure 3.1.

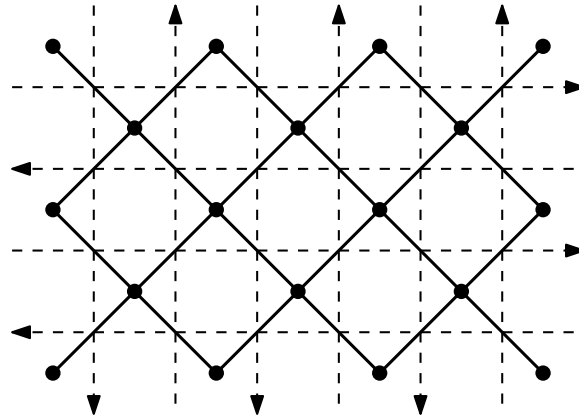


FIGURE 3.2. From the original (dashed) square lattice  $\mathbb{Z}^2$  one may construct a diagonal lattice  $\widehat{\mathbb{Z}}^2$ . In fact there are two such diagonal lattices, and this fact may be used to obtain some information about the power-law behaviour of the light ray when  $p = 1$ . The Manhattan orientations are not relevant to the usual Ehrenfest model, but are provided to facilitate the discussion of Manhattan pinball in Section 3.4.

Light is shone from the origin in a given compass direction, say north, and it is reflected off the surface of any mirror encountered. The problem is to decide whether or not the light ray is unbounded.

**Question 3.1.** *Let  $\theta(p)$  be the probability that the light ray is unbounded. For what values of  $p$  is it the case that  $\theta(p) > 0$ ?*

It is trivial that  $\theta(0) = 1$ . By considering bond percolation (with density  $p/2$ ) on the diagonal lattice of Figure 3.2, and using the fact that there is no percolation when  $p = 1$ , one obtains the less trivial fact that  $\theta(1) = 0$  (see [20]). Very little more is known rigorously about this question.

**3.3. Poisson mirrors.** Here is version of the wind/tree model in the two-dimensional continuum  $\mathbb{R}^2$ . Let  $\Pi$  be a rate-1 Poisson process in  $\mathbb{R}^2$ . Let  $\epsilon > 0$ , and let  $\mu$  be a probability measure on  $[0, \pi)$ . We possess an infinity of two-sided, plane mirrors of length  $\epsilon$ , and we centre one at each point in  $\Pi$ ; the inclination to the horizontal of each mirror is random with law  $\mu$ , and different mirrors have independent inclinations. Think of a mirror as being a randomly positioned, closed line segment of length  $\epsilon$ , and let  $M$  denote the union of these segments. We call  $\mu$

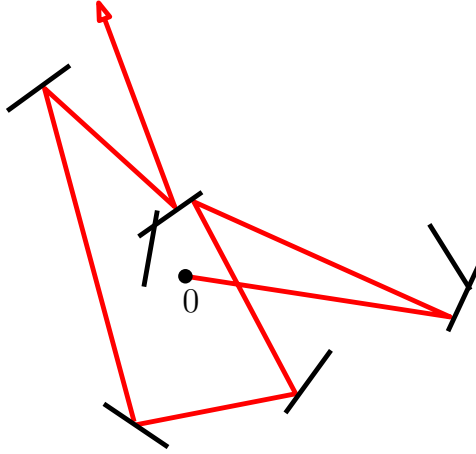


FIGURE 3.3. Light from the origin is reflected off the mirrors.

*degenerate* if it is concentrated on a single atom, and shall assume  $\mu$  is non-degenerate. See Figure 3.3.

Light is shone from the origin at an angle  $\alpha$  to the horizontal. Let  $I_\alpha$  be the indicator function of the event that the light ray is unbounded, and let  $\theta_\mu(\epsilon) = \mathbb{E}_\mu(I_\alpha)$  be the associated probability (here and later, the subscript  $\mu$  keeps track of the choice of  $\mu$ ). Some convention is adopted for the zero-probability event that the light strikes an intersection of two or more mirrors.

We may assume that the origin  $0$  does not lie in  $M$ . Let  $C$  be the component of  $\mathbb{R}^2 \setminus M$  containing  $0$ , and let  $\{0 \leftrightarrow \infty\}$  be the event that  $C$  is unbounded. It is a standard result of so-called needle percolation (see [49]) that there exists  $\epsilon_c = \epsilon_c(\mu) \in (0, \infty)$  such that

$$\mathbb{P}_\mu(0 \leftrightarrow \infty) \begin{cases} > 0 & \text{when } \epsilon < \epsilon_c, \\ = 0 & \text{when } \epsilon > \epsilon_c. \end{cases}$$

Obviously, on the event that  $0 \leftrightarrow \infty$ , we have that  $I_\alpha = 0$  for all  $\alpha$ . Therefore,

$$\mathbb{P}_\mu(I_\alpha = 1 \text{ for some } \alpha) = 0 = \theta_\mu(\epsilon), \quad \epsilon > \epsilon_c.$$

The converse issue is much harder and largely open.

**Question 3.2.** *Suppose  $\mu$  is non-degenerate.*

- (a) *Does there exist  $\epsilon'_c = \epsilon'_c(\mu) > 0$  such that  $\theta_\mu(\epsilon) > 0$  for  $\epsilon < \epsilon'_c$ ?*
- (b) *Could it be that  $\epsilon'_c(\mu) = \epsilon_c(\mu)$ ?*
- (c) *In particular, what happens when  $\mu$  is the uniform measure on  $[0, \pi)$ ?*

Let  $\mathcal{Q}$  be the set of probability measures  $\mu$  that are non-degenerate and have support in the rational angles  $\pi\mathbb{Q}$ . Suppose  $\mu \in \mathcal{Q}$  and  $0 < \epsilon < \epsilon_c(\mu)$ . Harris [31] has shown the striking fact that,  $\mathbb{P}_\mu$ -a.s. on the event that  $0 \leftrightarrow \infty$ , we have that  $I_\alpha = 1$  for (Lebesgue) almost every  $\alpha$ .

This leads to a deterministic question. Let  $K$  be the set of mirror configurations for which  $0 \leftrightarrow \infty$  but  $I_\alpha = 0$  for all  $\alpha$ .

**Question 3.3.** *Is  $K$  non-empty?*

Harris's theorem states in effect that  $\mathbb{P}_\mu(K) = 0$  when  $\mu \in \mathcal{Q}$  and  $\epsilon \neq \epsilon_c(\mu)$ . Question 3.2(c) hints at the possibility that  $\mathbb{P}_\mu(K) = 0$  when  $\mu$  is the uniform measure on  $[0, \pi)$  and  $\epsilon < \epsilon_c(\mu)$ .

Here is a final question concerning diffusivity. Let  $\mu \in \mathcal{Q}$ , and denote by  $X_\alpha(t)$  the position at time  $t$  of the light ray that leaves the origin at angle  $\alpha$ .

**Question 3.4.** *Is it the case that, on the event  $I_\alpha$ ,  $X_\alpha(\cdot)$  is diffusive? That is, the limit  $\sigma^2 := \lim_{t \rightarrow \infty} t^{-1} \text{var}(X_\alpha(t))$  exists in  $(0, \infty)$ , and, when normalized,  $X_\alpha(t)$  is asymptotically normally distributed.*

Related work on Lorentz models in the so-called Boltzmann–Grad limit may be found in [39, 40].

**3.4. Manhattan pinball.** Here is a variant of the Ehrenfest model motivated by a problem of quantum localization, [13, Sec. 4.2], [50, p. 238]. Draw  $\mathbb{Z}^2$  and the diagonal lattice  $\widehat{\mathbb{Z}}^2$  as in Figure 3.2; each edge of  $\mathbb{Z}^2$  receives its Manhattan orientation as indicated in the figure. Consider bond percolation with density  $q$  on the diagonal lattice. Along each open edge of  $\widehat{\mathbb{Z}}^2$  we place a two-sided plane mirror. Light is shone from the origin along a given one of the two admissible directions, and it is reflected by any mirror that it encounters (such reflections are automatically consistent with the Manhattan orientations). Let  $\theta(q)$  be the probability that the light ray is unbounded.

**Question 3.5.** *Could it be that  $\theta(q) = 0$  for all  $q > 0$ ?*

It follows as in Section 3.2 that  $\theta(q) = 0$  for  $q \geq \frac{1}{2}$ , and it has been proved by Li in [37] that there exists  $\epsilon > 0$  such that  $\theta(q) = 0$  when  $q > \frac{1}{2} - \epsilon$ . The proof uses the method of enhancements; see Section 6.

## 4. TWO STOCHASTIC INEQUALITIES

**4.1. Bunkbed inequality.** The mysterious ‘bunkbed’ inequality was posed by Kasteleyn around 1985 (see [9, Rem. 5], and also [28]). Of its various flavours, we select the following. Let  $G = (V, E)$  be a finite

simple graph. From  $G$  we construct two copies denoted  $G_1 = (V_1, E_1)$  and  $G_2 = (V_2, E_2)$ . For  $v \in V$  we write  $v_i$  for the copy of  $v$  lying in  $V_i$ . We now stick  $G_1$  and  $G_2$  together by adding edges  $\langle v_1, v_2 \rangle$  for  $v \in V$ . This new graph is denoted  $\tilde{G}$ , and it may be considered as the product graph  $G \times K_2$  where  $K_2$  is the complete graph on two vertices (that is, an edge). We think of the  $G_i$  as being ‘horizontal’ and the extra edges as being ‘vertical’.

Each edge of  $\tilde{G}$  is declared *open* with probability  $p$ , independently of the states of other edges. Write  $\mathbb{P}_p$  for the appropriate product measure. For two vertices  $u_i, v_j$  of  $\tilde{G}$ , we write  $\{u_i \leftrightarrow v_j\}$  for the event that there exists a  $u_i/v_j$  path using only open edges.

**Conjecture 4.1.** *For  $u, v \in V$ , we have*

$$\mathbb{P}_p(u_1 \leftrightarrow v_1) \geq \mathbb{P}_p(u_1 \leftrightarrow v_2).$$

There is uncertainty over whether this was the exact conjecture of Kasteleyn. For example, it is suggested in [33] (and perhaps elsewhere) that Kasteleyn may have made the stronger conjecture that the inequality holds even after conditioning on the set  $T$  of open vertical edges.

Some special cases of the bunkbed conjecture have been proved (see the references in [33]), but the general question remains open.

**4.2. Negative correlation.** Our next problem is quite longstanding (see [47]) and remains mysterious. In a nutshell it is to prove that the uniform random forest measure (USF) has a property of negative association.

Let  $G = (V, E)$  be a finite graph which, for simplicity, we assume has neither loops nor multiple edges. A subset  $F \subseteq E$  is called a *forest* if  $(V, F)$  has no cycles. Let  $\mathcal{F}$  be the set of all forests in  $G$  and let  $\Phi$  be a random forest chosen uniformly from  $\mathcal{F}$ . We call  $\Phi$  *edge-negatively associated* if

$$(4.1) \quad \mathbb{P}(e, f \in \Phi) \leq \mathbb{P}(e \in \Phi)\mathbb{P}(f \in \Phi), \quad e, f \in E, e \neq f.$$

**Conjecture 4.2.** *For all graphs  $G$ , the random forest  $\Phi$  is edge-negatively associated.*

One may formulate various forms of negative dependence, amongst which (4.1) is quite weak. One may conjecture that  $\Phi$  has a stronger variety of such dependence. Further discussion may be found in [22, Sec. 3.9] and [47].

Experimental evidence for Conjecture 4.2 is quite strong. A similar conjecture may be made for uniform measure on the set of  $F \subseteq E$  such



that  $(V, F)$  is connected (abbreviated to UCS). In contrast, uniform spanning tree (UST) is well understood via the Kirchhoff theory of electrical networks, and further by [16]. USF, UCS, and UST are special cases of the so-called random-cluster measure with cluster-weighting factor  $q$  satisfying  $q < 1$  (see [22]).

In recent work, [5, 6], the percolative properties of the weighted random forest (or ‘arboreal gas’) on  $\mathbb{Z}^d$  have been explored. It turns that there is a phase transition if and only if  $d \geq 3$ .

## 5. RANDOMLY ORIENTED SQUARE LATTICE

The following percolation-type problem remains open. Consider the square lattice  $\mathbb{Z}^2$  and let  $p \in [0, 1]$ . Each horizontal edge is oriented rightward with probability  $p$ , and otherwise leftward. Each vertical edge is oriented upward with probability  $p$ , and otherwise downward. Write  $\vec{\mathbb{Z}}^2$  for the ensuing randomly oriented network.

Let  $\theta(p)$  denote the probability that the origin 0 is the endpoint of an infinite path of  $\vec{\mathbb{Z}}^2$  that is oriented away from 0. The challenge is to determine for which  $p$  it is the case that  $\theta(p) > 0$ . It is elementary that  $\theta(0) = 1$ , and that  $\theta(p) = \theta(1 - p)$ . It is less obvious that  $\theta(\frac{1}{2}) = 0$  (see [20, 1st edn]), which is proved via a coupling with bond percolation. By a comparison with oriented percolation, we have that  $\theta(p) > 0$  if  $p > \vec{p}_c$ , where  $\vec{p}_c$  is the critical point of oriented percolation on  $\mathbb{Z}^2$ ; it is not difficult to deduce by the enhancement method of [2] (see Section 6) that there exists  $p' \in (\frac{1}{2}, \vec{p}_c)$  such that  $\theta(p) > 0$  when  $p > p'$ . It is believed that  $\vec{p}_c \sim 0.64$ , and proved that  $\vec{p}_c < 0.6735$ .

**Question 5.1.** *Is it the case that  $\theta(p) > 0$  for  $p \neq \frac{1}{2}$ .*

It is shown in [21] that, for all  $p$ ,  $\vec{\mathbb{Z}}^2$  is either critical or supercritical in the following sense: if any small positive density of oriented edges is added at random, then there is a strictly positive probability that the origin is the endpoint of an infinite self-avoiding oriented path in the resulting graph.

## 6. ENHANCEMENTS

The technique of so-called enhancements was introduced in [2] to answer the question “when is the critical point of a disordered system a *strictly* monotone function of the parameters of the system?” There have been numerous applications of this technique, including to percolation [1, 12, 14, 18, 26, 27, 32, 43], the Ising/Potts and random-cluster models [10, 19], as well as the Lorentz gas [37] and the minesweeper game [44].

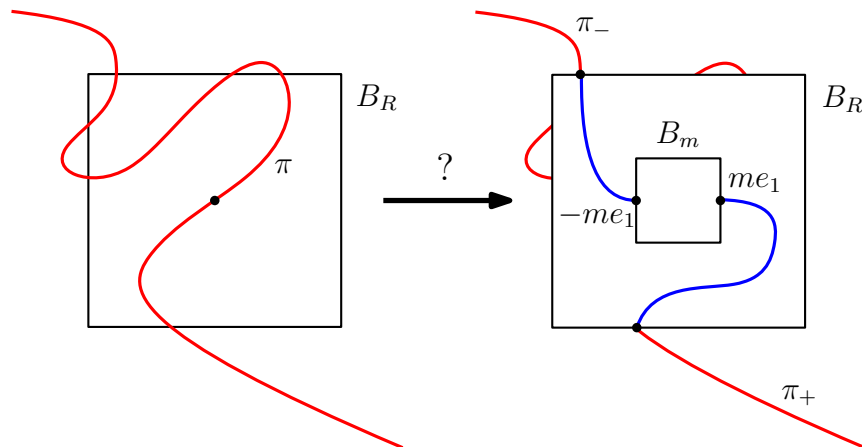


FIGURE 6.1. By making changes only within  $B_R$ , can one replace a non-self-touching path  $\pi$  by a non-touching pair  $\pi_-, \pi_+$  of paths?

The approach of [2] was to describe a conceptual technique with wide applicability, illustrated by applications. Some of the graph-theoretic details of the application to site percolation were omitted, and this has attracted the attention of the authors of [4] who have identified a difficulty in the special case of site percolation on  $\mathbb{Z}^d$  with  $d \geq 3$ . The corresponding step for *bond* percolation was treated in adequate detail in [20, Chap. 3].

Let  $d \geq 2$  and let  $\mathbb{Z}^d$  be the set of  $d$ -vectors  $v = (v_1, v_2, \dots, v_d)$  of integers. We turn  $\mathbb{Z}^d$  into a graph by adding an edge  $\langle v, w \rangle$  between any pair  $v, w \in \mathbb{Z}^d$  satisfying  $\sum_{i=1}^d |v_i - w_i| = 1$ .

A *configuration* is an element  $\omega$  of the space  $\Omega := \{0, 1\}^{\mathbb{Z}^d}$ . The site percolation model is obtained by choosing  $\omega$  according to the product measure  $\mathbb{P}_p$  on  $\Omega$ , but probability plays no direct role in the following. We say that a configuration  $\omega \in \Omega$  *contains* a vertex  $v \in \mathbb{Z}^d$  if  $\omega(v) = 1$ .

A path  $\pi$  of  $\mathbb{Z}^d$  is an alternating sequence  $\dots, x_0, e_0, x_1, e_1, \dots, e_{n-1}, x_n, \dots$  of vertices  $x_i$  and edges  $e_i$  such that  $e_i = \langle x_i, x_{i+1} \rangle$  for all  $i$ . The path  $\pi$  is called *non-self-touching* if  $\delta(x_i, x_j) \geq 2$  whenever  $j - i \geq 2$ , where  $\delta(u, v)$  denotes the number of edges in the shortest path from  $u$  to  $v$ . Two paths with respective vertex-sequences  $(x_i)$  and  $(y_j)$  are called *non-touching* if each is non-self-touching and, in addition,  $\delta(x_i, y_j) \geq 2$  for all  $i, j$ .

Let  $e_1 = (1, 0, \dots, 0)$  be the unit vector with 1 in the first position, and let  $B_m = [-m, m]^d \cap \mathbb{Z}^d$ .

**Question 6.1.** *Let  $m \geq 1$  and  $d \geq 2$ . Does there exist  $R > m$  such that the following holds? Let  $\omega \in \Omega$  be a configuration that contains a doubly infinite, non-self-touching path  $\pi$  containing the origin 0. There exists a configuration  $\omega'$  that agrees with  $\omega$  off  $B_R$  such that:*

- (a)  $\omega'$  contains a singly infinite path  $\pi_-$  with initial vertex  $-me_1$  and no other vertex belonging to  $B_m$ ,
- (b)  $\omega'$  contains a singly infinite path  $\pi_+$  with initial vertex  $me_1$  and no other vertex belonging to  $B_m$ ,
- (c) the pair  $(\pi_-, \pi_+)$  is non-touching.

This is illustrated in Figure 6.1. The motivation for the question is as follows. Suppose the vertex 0 is pivotal for a certain event  $E$ , and it is proposed to apply an enhancement that has the effect of adding certain open connections within the inner box  $B_m$ . Then, by changing the configuration within some bounded box  $B_R$ , one may obtain a new configuration in which the enhancement is itself pivotal for  $E$ . By Russo's formula, this implies the comparability of the two partial derivatives of  $\mathbb{P}_{p,\eta}(A)$ , where  $\eta$  is the probability that any given enhancement is 'activated'. The strict monotonicity of the critical point follows on integration (see [20, Sec. 3.3] for further details).

It was implicit in the proof of [2, Thm 1] that Question 6.1 has an affirmative answer for  $d \geq 2$ . The proof is straightforward when  $d = 2$ . A positive answer is proved in [4] when  $d = 3$ , and the question seems to be open for  $d \geq 4$ .

A closely related issue arises in the so-called Burton–Keane [11] proof of the uniqueness of the infinite percolation cluster. Path-surgery of the above type may in fact be avoided in the latter context, see [8, 41].

## 7. TWO STOCHASTIC EPIDEMICS

The recent pandemic has inspired a number of mathematical problems, including the following stochastic model (see [24]). Particles are placed at time 0 at the points of a rate-1 Poisson process in  $\mathbb{R}^d$ , where  $d \geq 1$ . Each particle diffuses around  $\mathbb{R}^d$  according to a Brownian motion, independently of other particles. At any given time, each particle is in one of the states S (susceptible), I (infected), R (removed/dead).

At time 0 there exists a unique particle in state I, and all others are in state S. The infection/removal rules are as follows.

- (a) If an infected particle comes within distance 1 of a susceptible particle, the latter particle is infected.
- (b) An infected particle remains infected for a period of time having the exponential distribution with parameter  $\alpha$ , and is then removed.

We call this the ‘diffusion model’.

*Survival* is said to occur if, with a strictly positive probability, infinitely many particles are ultimately infected. It is proved in [24] that, when  $d \geq 1$  and  $\alpha$  is sufficiently large, survival does not occur. The following two questions (amongst others) are left open.

**Question 7.1.**

- (i) *When  $d = 1$ , could it be that there is no survival for any  $\alpha > 0$ ?*
- (ii) *When  $d \geq 2$ , does survival occur for sufficiently small  $\alpha > 0$ ?*

In a variant of this problem termed the ‘delayed diffusion model’, a much fuller picture is known. Suppose, instead of the above, a particle moves only when it is infected; susceptible particles are stationary. The answers to Question 7.1(i, ii) are then no and yes, respectively, and indeed (when  $d \geq 2$ ) there exists a critical value  $\alpha_c(d) \in (0, \infty)$  of  $\alpha$  marking the onset of survival. The key difference between the two systems is that the latter model has a property of monotonicity that is lacking in the former.

The delayed diffusion model is a continuous space/time cousin of the discrete-time ‘frog model’ of [3] (see also [48]), with the addition with removal. Further relevant references may be found in [24].

## 8. ACKNOWLEDGEMENTS

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